# Localization of peaks for high-order equations 

Frédéric Robert<br>Institut Elie Cartan, Université de Lorraine (Nancy)

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The equation
Let ( $M, g$ ) be a compact Riemmanian manifold of dimension $n \geq 2$, and take $k \in \mathbb{N}$ such that $n>2 k \geq 2$. We are interested in functions $u \in C^{2 k}(M)$ that are solutions to

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P u=|u|^{2^{\star}-2} u \text { in } M \text { with } 2^{\star}:=\frac{2 n}{n-2 k}
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Here, $\Delta_{g}=-\operatorname{div}_{g} \nabla$. Such a PDE arises in conformal geometry:

- $\mathbf{k}=\mathbf{1}$, the scalar curvature equation is

$$
\Delta_{g} u+\frac{n-2}{4(n-1)} R_{g} u=\frac{n-2}{4(n-1)} R_{\tilde{g}} u^{\frac{n+2}{n-2}}, u>0
$$

where $R_{g}$ (resp. $R_{\tilde{g}}$ ) is the scalar curvature of $g$ (resp. $\tilde{g}=u^{\frac{4}{n-2}} g$ ).

- $\mathbf{k}=\mathbf{2}$, the Paneitz operator connects Branson's $Q$-curvatures in a conformal class too:

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\Delta_{g}^{2} u+\ldots=Q_{\tilde{g}} u^{\frac{n+4}{n-4}}
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- More generally, for any $\mathbf{k} \geq \mathbf{1}$, there is the conformal GJMS operator $P_{g}$ and a notion of $Q$-curvature
These operators are conformally invariant in the following sense: if $\tilde{g}=u^{\frac{4}{n-2 k}} g$, then

$$
P_{\tilde{g}} \varphi=u^{-\left(2^{\star}-1\right)} P_{g}(u \varphi) \text { for all } \varphi \in C^{\infty}(M)
$$

The invariance/instability of the equation
When $(\mathbf{M}, \mathbf{g})=\left(\mathbb{R}^{\mathbf{n}}, \xi\right)$ (which is not compact...) the model is

$$
\Delta_{\text {eucl }}^{k} U=U^{2^{\star}-1}, U>0 \text { in } \mathbb{R}^{n} .
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\Delta_{\xi}^{k} U_{\mu, x_{0}}=U_{\mu, x_{0}}^{2^{\star}-1}, U_{\mu, x_{0}}>0 \text { in } \mathbb{R}^{n} .
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This invariance generates an intrinsic dynamic of the equation.
For instance, you can take

$$
U(x):=\alpha_{n, k}\left(\frac{1}{1+|x|^{2}}\right)^{\frac{n-2 k}{2}} \quad U^{\frac{4}{n-2 k}}=\alpha_{n, k}^{\prime}\left(\frac{1}{1+|x|^{2}}\right)^{2} \Rightarrow \text { round sphere }
$$

so that

$$
U_{\mu, x_{0}}(x):=\alpha_{n, k}\left(\frac{\mu}{\mu^{2}+\left|x-x_{0}\right|^{2}}\right)^{\frac{n-2 k}{2}}
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so that

$$
\lim _{\mu \rightarrow 0} U_{\mu, x_{0}}\left(x_{0}\right)=+\infty \text { and } \lim _{\mu \rightarrow 0} U_{\mu, x_{0}}(x)=0 \text { for all } x \neq x_{0}
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Figure: $\lim _{\mu \rightarrow 0} U_{\mu, x_{0}}\left(x_{0}\right)=+\infty$ and $\lim _{\mu \rightarrow 0} U_{\mu, x_{0}}(x)=0$ for all $x \neq x_{0}$

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And they are going to be our model to describe instability But there can be other types of peaks.

## Definition (Exponential chart)

A smooth exponential chart ex̃p around $p_{0} \in M$ is a function

$$
\begin{array}{cccc}
\operatorname{exx}_{p}: & \mathbb{R}^{n} & \rightarrow & M \\
& \left(X^{1}, \ldots, X^{n}\right) & \mapsto & \exp _{p}\left(\sum_{i} X^{i} E_{i}(p)\right)
\end{array}
$$

where $\exp _{p}: T_{p} M \rightarrow M$ is the usual exponential map and $\left(E_{i}(p)\right)_{i=1, \ldots, n}$ is a smooth orthonormal basis of $T_{p} M, p$ close to $p_{0}$.

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## Definition (Peak)

We say that a family $B=\left(B_{\alpha}\right)_{\alpha} \in H_{k}^{2}(M)$ is a Peak centered at $\left(x_{\alpha}\right)_{\alpha} \in M$ with radius $\left(\mu_{\alpha}\right)_{\alpha} \rightarrow 0$ if there exists $U \in D_{k}^{2}\left(\mathbb{R}^{n}\right), U \not \equiv 0$, and an exponential chart e $\tilde{x} p$ around $x_{0}:=\lim _{\alpha \rightarrow 0} x_{\alpha}$ and a cutoff function $\left(\eta_{\alpha}\right)$ such that

$$
\begin{equation*}
B_{\alpha}(x)=\eta(x) \mu_{\alpha}^{-\frac{n-2 k}{2}} U\left(\frac{e \tilde{x} p_{x_{\alpha}}^{-1}(x)}{\mu_{\alpha}}\right)+o(1) \text { in } H_{k}^{2}(M) . \tag{1}
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There are examples of solutions to our PDE blowing-up like peaks (Pistoia et al, Vétois et al, Casteras-Bakri). We prove here that any blowing-up solutions behave like a peak and has a precise localization:

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## Theorem (R., 2023)

Consider a family $\left(u_{\alpha}\right)_{\alpha} \in C^{2 k}(M)$ such that

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\begin{aligned}
& \Delta_{g}^{k} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(A_{\alpha} \nabla^{k-1} u_{\alpha}\right)+\text { lot }=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M, \text { for all } \alpha>0 . \\
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\text { •If } n>2 k+2 \text { (similar for } n=2 k+2) \text {, then } \\
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Moreover,

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\left|u_{\alpha}(x)\right| \leq C\left(\frac{\mu_{\alpha}}{\mu_{\alpha}^{2}+d_{g}\left(x, x_{\alpha}\right)^{2}}\right)^{\frac{n-2 k}{2}}
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- It is $=0$ when $U$ is radial


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$\Rightarrow$ the second term measures the "distance" of the limiting op. to the geometric op.

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\text { Weyl }_{g} \otimes B:=\left(\text { Weylg }_{g}\left(x_{0}\right)\right)_{i \alpha j \beta} \int_{\mathbb{R}^{n}} X^{\alpha} X^{\beta} \partial_{i j} \Delta_{\text {eucl }}^{k-1} U\left(\frac{n-2 k}{2} U+X^{\prime} \partial_{l} U\right) d X
$$

- It is independent of the choice of $U$ in the definition of the peak,
- It is $=0$ when Weyl vanishes at $x_{0}$
- It is $=0$ when $U$ is radial
- It is $=0$ when $u_{\alpha}>0$ (since then, $U>0$ and is then radial wrt a point)


## A closer look at the case $n>2 k+2$

If $u_{\alpha}=B_{\alpha}+o(1)$ for a bubble $B=\left(B_{\alpha}\right)_{\alpha}$ where $P_{\alpha} u_{\alpha}=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha}$, then

$$
\text { Weyl }_{g} \otimes B+\int_{\mathbb{R}^{n}}\left(A_{\infty}-A_{G J M S}\right)_{x_{0}}\left(\nabla^{k-1} U, \nabla^{k-1} U\right) d X=0
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Here

$$
P_{\alpha} \rightarrow P_{\infty}=\Delta_{g}^{k}+(-1)^{k-1} \nabla^{k-1}\left(A_{\infty} \nabla^{k-1}\right)+\operatorname{lot}
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$\Rightarrow$ Weyl ${ }_{g} \otimes B$ arises only when dealing with sign-changing $u_{\alpha}$ in the non-Icf setting

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$$
\int_{\mathbb{R}^{n}}|U|^{2^{\star}-2} U d X=0 \Leftrightarrow U(x)=o\left(|x|^{2 k-n}\right) \text { as }|x| \rightarrow \infty
$$

This is possible only for some sign-changing $U$, but not all of them.

## Theorem (R., 2023)

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\left\{\begin{array}{c}
P_{\alpha} u_{\alpha}=\left|u_{\alpha}\right|^{\star}-2 u_{\alpha} \text { in } M \\
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The key is to get the pointwise control.

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- Here: we want a method in analysis that does not require geometric assumptions or sign assumptions... because it simply more natural.
- The main difficulty: how to bypass the maximum principle?

Proof. Step 1: rescaling

$$
\Delta_{\xi}^{k} u_{\alpha}+\ldots=u_{\alpha}^{2^{\star}-1}, u_{\alpha}>0 \text { (for simplification) }
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- Elliptic regularity:

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Objective: We want ( $\star$ ) on all the manifold $M$

Proof. Step 2: the final argument in 1 page!

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- Let $G_{\alpha}$ be the Green's function for $P_{\alpha}-V_{\alpha}$
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$$
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- Write $P_{\alpha} u_{\alpha}=u_{\alpha}^{2^{\star}-1}$ as a linear problem

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SO, we must get a reasonable control for the Green's function of

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with a potential that is blowing-up.Fortunately, this is a very particular potential.

Intermission: $V_{\alpha}=u_{\alpha}^{2^{\star}-2}$ is a Hardy-potential
Let us go back to the invariance. A paramount quantity here is

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w_{\alpha}(x):=d_{g}\left(x, x_{\alpha}\right)^{\frac{n-2 k}{2}} u_{\alpha}(x)
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\lim _{\mathbf{R} \rightarrow+\infty} \lim _{\alpha \rightarrow+\infty} \sup _{\mathbf{M}-\mathbf{B}\left(\mathbf{x}_{\alpha}, \mathbf{R} \mu_{\alpha}\right)} \mathbf{d}_{\mathbf{g}}\left(\mathbf{x}, \mathbf{x}_{\alpha}\right)^{\frac{\mathbf{n}-2 \mathbf{k}}{2}}\left|\mathbf{u}_{\alpha}(\mathbf{x})\right|=\mathbf{0}
$$

So $\forall \epsilon>0, \exists R_{\epsilon}>0$ such that $\left|V_{\alpha}(x)\right| \leq \frac{\epsilon}{d_{g}\left(x, x_{\alpha}\right)^{2 k}}$ for all $d_{g}\left(x, x_{\alpha}\right)>R_{\epsilon} \mu_{\alpha}$

## Back to the proof: What we consider now

Our equation rewrites

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\left(P_{\alpha}-V_{\alpha}\right) u_{\alpha}=0
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and

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In order to perform the argument that failed earlier, we need to get a pointwise control on $G_{\alpha}$, the Green's function of $P_{\alpha}-V_{\alpha}$ with the property above.

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- When $x, y \in M$ are far from the singularity $x_{\alpha}$, then

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\left|G_{\alpha}(x, y)\right| \leq C d_{g}(x, y)^{2 k-n} \Rightarrow \operatorname{Good}
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- For general $x, y$ : a mix of these two cases.

We get a sharp control of the Green's function and of its derivatives

At the end of the day, we have proved that

## Theorem (R., 2022)

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\left\{\begin{array}{c}
P_{\alpha} u_{\alpha}=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M \\
+ \text { Blow-up } \\
+ \text { Minimal energy }
\end{array}\right\} \Rightarrow\left|u_{\alpha}(x)\right| \leq C\left(\frac{\mu_{\alpha}}{\mu_{\alpha}^{2}+d_{g}\left(x, x_{\alpha}\right)^{2}}\right)^{\frac{n-2 k}{2}}
$$

where $\mu_{\alpha}^{-\frac{n-2 k}{2}}=\left|u_{\alpha}\left(x_{\alpha}\right)\right|=\sup _{M}\left|u_{\alpha}\right|$.

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+ \text { Blow-up } \\
+ \text { Minimal energy }
\end{array}\right\} \Rightarrow\left|u_{\alpha}(x)\right| \leq C\left(\frac{\mu_{\alpha}}{\mu_{\alpha}^{2}+d_{g}\left(x, x_{\alpha}\right)^{2}}\right)^{\frac{n-2 k}{2}}
$$

where $\mu_{\alpha}^{-\frac{n-2 k}{2}}=\left|u_{\alpha}\left(x_{\alpha}\right)\right|=\sup _{M}\left|u_{\alpha}\right|$.

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
\int_{\Omega}\left(x^{i} \partial_{i} v+\frac{n-2 k}{2} v\right)\left(\Delta_{\xi}^{k} v-|v|^{2^{\star}-2} v\right) d x=\int_{\partial \Omega} \ldots
$$

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
\int_{\Omega}\left(x^{i} \partial_{i} v+\frac{n-2 k}{2} v\right)\left(\Delta_{\xi}^{k} v-|v|^{2^{\star}-2} v\right) d x=\int_{\partial \Omega} \cdots
$$

Our equation is

$$
\Delta_{g}^{k} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(A_{\alpha} \nabla^{k-1} u_{\alpha}\right)+\text { lot }=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M
$$

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$$
\Delta_{g}^{k} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(A_{\alpha} \nabla^{k-1} u_{\alpha}\right)+l o t=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M
$$

we write it as

$$
P_{g} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(\left(A_{\alpha}-A_{G J M S}\right) \nabla^{k-1} u_{\alpha}\right)+\text { lot }=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M
$$

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
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$$

we can change the metric in metric $\tilde{g}$ conformal to $g$ becomes "almost flat", that is $\operatorname{Ric}_{\tilde{g}}\left(x_{\alpha}\right)=0$.

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
\int_{\Omega}\left(x^{i} \partial_{i} v+\frac{n-2 k}{2} v\right)\left(\Delta_{\xi}^{k} v-|v|^{2^{\star}-2} v\right) d x=\int_{\partial \Omega} \cdots
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$$

we can change the metric in metric $\tilde{g}$ conformal to $g$ becomes "almost flat", that is $\operatorname{Ri} \tilde{\varepsilon}_{\tilde{g}}\left(x_{\alpha}\right)=0$. We then write

$$
\Delta_{\tilde{g}}^{k} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(\left(A_{\alpha}-A_{G J M S}\right) \nabla^{k-1} u_{\alpha}\right)+l o t=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M
$$

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
\int_{\Omega}\left(x^{i} \partial_{i} v+\frac{n-2 k}{2} v\right)\left(\Delta_{\xi}^{k} v-|v|^{2^{\star}-2} v\right) d x=\int_{\partial \Omega} \cdots
$$

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$$

In the Pohozaev identity, we then get

$$
\int_{\Omega}\left(x^{i} \partial_{i} u_{\alpha}+\frac{n-2 k}{2} u_{\alpha}\right)\left(\left(\Delta_{\xi}^{k}-\Delta_{\tilde{g}}^{k}\right) u_{\alpha}-(-1)^{k-1} \nabla^{k-1}\left(\left(A_{\alpha}-A_{G J M S}\right) \nabla^{k-1} u_{\alpha}\right)\right) d x=\ldots
$$

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
\int_{\Omega}\left(x^{i} \partial_{i} v+\frac{n-2 k}{2} v\right)\left(\Delta_{\xi}^{k} v-|v|^{2^{\star}-2} v\right) d x=\int_{\partial \Omega} \cdots
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$$

we can change the metric in metric $\tilde{g}$ conformal to $g$ becomes "almost flat", that is $\operatorname{Ri} \tilde{\tilde{g}}_{\tilde{g}}\left(x_{\alpha}\right)=0$. We then write

$$
\Delta_{\tilde{\mathrm{g}}}^{k} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(\left(A_{\alpha}-A_{G J M S}\right) \nabla^{k-1} u_{\alpha}\right)+l o t=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M
$$

In the Pohozaev identity, we then get
$\int_{\Omega} T\left(u_{\alpha}\right)(\underbrace{\left(\Delta_{\xi}^{k}-\Delta_{\tilde{g}}^{k}\right) u_{\alpha}}_{\text {measures } \tilde{g}-\xi}-(-1)^{k-1} \nabla^{k-1}(\underbrace{\left(A_{\alpha}-A_{G J M S}\right)}_{\text {distance from the conf.op. }} \nabla^{k-1} u_{\alpha})) d x=\ldots$
where $T\left(u_{\alpha}\right):=x^{i} \partial_{i} u_{\alpha}+\frac{n-2 k}{2} u_{\alpha}$.

For any function $v: \mathbb{R}^{n} \rightarrow \mathbb{R}$ and $\Omega \subset \mathbb{R}^{n}$, we have that

$$
\int_{\Omega}\left(x^{i} \partial_{i} v+\frac{n-2 k}{2} v\right)\left(\Delta_{\xi}^{k} v-|v|^{2^{\star}-2} v\right) d x=\int_{\partial \Omega} \cdots
$$

Our equation is

$$
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$$

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$$
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$$

we can change the metric in metric $\tilde{g}$ conformal to $g$ becomes "almost flat", that is $\operatorname{Ri} \tilde{\tilde{g}}_{\tilde{g}}\left(x_{\alpha}\right)=0$. We then write

$$
\Delta_{\tilde{\mathrm{g}}}^{k} u_{\alpha}+(-1)^{k-1} \nabla^{k-1}\left(\left(A_{\alpha}-A_{G J M S}\right) \nabla^{k-1} u_{\alpha}\right)+l o t=\left|u_{\alpha}\right|^{2^{\star}-2} u_{\alpha} \text { in } M
$$

In the Pohozaev identity, we then get

$$
\int_{\Omega} T\left(u_{\alpha}\right)(\underbrace{\left(\Delta_{\xi}^{k}-\Delta_{\tilde{g}}^{k}\right) u_{\alpha}}_{\text {measures } \tilde{g}-\xi}-(-1)^{k-1} \nabla^{k-1}(\underbrace{\left(A_{\alpha}-A_{G J M S}\right)}_{\text {distance from the conf.op. }} \nabla^{k-1} u_{\alpha})) d x=\ldots
$$

where $T\left(u_{\alpha}\right):=x^{i} \partial_{i} u_{\alpha}+\frac{n-2 k}{2} u_{\alpha}$. When $n>2 k+2$, we get

$$
\text { Weyl }_{g} \otimes B+\int_{\mathbb{R}^{n}}\left(A_{\infty}-A_{G J M S}\right)_{x_{0}}\left(\nabla^{k-1} U, \nabla^{k-1} U\right) d X=0
$$

On ( $M, g$ ) of dimension $n \geq 5$, see Hebey, there exists $B>0$ such that the following Sobolev inequality holds:

$$
\left(\int_{M}|u|^{\frac{2 n}{n-4}} d v_{g}\right)^{\frac{n-4}{n}} \leq K_{4}(n)\left(\int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+B\|u\|_{H_{1}^{2}}^{2}\right) \text { for all } u \in H_{2}^{2}(M) . \quad\left(I_{B}\right)
$$

where $K_{4}(n)$ is the optimal Euclidean constant.

On $(M, g)$ of dimension $n \geq 5$, see Hebey, there exists $B>0$ such that the following Sobolev inequality holds:

$$
\left(\int_{M}|u|^{\frac{2 n}{n-4}} d v_{g}\right)^{\frac{n-4}{n}} \leq K_{4}(n)\left(\int_{M}\left(\Delta_{g} u\right)^{2} d v_{g}+B\|u\|_{H_{1}^{2}}^{2}\right) \text { for all } u \in H_{2}^{2}(M)
$$

where $K_{4}(n)$ is the optimal Euclidean constant.Let $B_{0}(g)$ be the smallest number $B$ such that this inequality holds for all $u \in H_{2}^{2}(M)$.

## Theorem

Assume that $n \geq 6$. Then if there is no nontrivial extremal for $\left(I_{B_{0}(g)}\right)$, then

$$
B_{0}(g)=\frac{3 n^{2}-6 n-12}{6 n(n-1)} \max _{x \in M} R_{g}(x)
$$

$$
\begin{gathered}
\operatorname{Pag} u=\Delta_{g}^{2} u-\operatorname{div}_{g}\left[\left(a_{n} S_{g} g+b_{n} R i c_{g}\right)^{\#} d u\right]+\frac{n-4}{2} Q_{g} u, \\
a_{n}=\frac{(n-2)^{2}+4}{2(n-1)(n-2)}, \quad b_{n}=-\frac{4}{n-2}, \\
Q_{g}^{n}=\frac{1}{2(n-1)} \Delta_{g} R_{g}+\frac{n^{3}-4 n^{2}+16 n-16}{8(n-1)^{2}(n-2)^{2}} R_{g}^{2}-\frac{2}{(n-2)^{2}}\left|R i c_{g}\right|_{g}^{2} .
\end{gathered}
$$

